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# CODES FOR CHAPTERS 2 AND 3

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## Chapters 2 and 3

### **fresnel.f, fresnel.cpp, fresnel.m**

Subroutine `fresnel(n,k,th,rhos,rhop,rho)` calculates Fresnel reflectances from equation (2.113).

Input: `n` ( $= n$ ) and `k` ( $= k$ ) are real and imaginary parts of the complex index of refraction, and `th` ( $= \theta$ ) is the off-normal angle of incidence (in radians).

Output: `rhos` ( $= \rho_{\perp}$ ) and `rhop` ( $= \rho_{\parallel}$ ) are perpendicular and parallel-polarized reflectance, respectively, while `rho` ( $= \rho$ ) is the unpolarized reflectance.

## Chapter 3

### **emdiel.f90, emdiel.cpp, emdiel.m**

Function `emdiel(n)` calculates the unpolarized, spectral, hemispherical emissivity of an optical surface of a dielectric material from equation (3.82).

Input: `n` ( $= n$ ) refractive index of dielectric.

### **emmet.f90, emmet.cpp, emmet.m**

Function `emmet(n,k)` calculates the unpolarized, spectral, hemispherical emissivity of an optical surface of a metallic material from equation (3.77).

Input: `n` ( $= n$ ) and `k` ( $= k$ ) are the real and imaginary parts of the metal's complex index of refraction.

### **callemdiel.f90, callemdiel.cpp, emmet.m, callemdiel.exe**

Program `callemdiel` is a stand-alone front end for function `emdiel`, prompting for input (refractive index  $n$ ) and returning the unpolarized, spectral, hemispherical as well as normal emissivities.

### **callemmet.f90, callemmet.cpp, callemmet.m, callemmet.exe**

Program `callemmet` is a stand-alone front end for function `emmet`, prompting for input (complex index of refraction  $n, k$ ) and returning the unpolarized, spectral, hemispherical as well as normal emissivities.

### **dirreflec.f, dirreflec.cpp, dirreflec.m, dirreflec.exe**

Program `dirreflec` is a stand-alone front end for subroutine `fresnel`, calculating reflectivities for various incidence angles. The user is prompted to input the complex index of refraction,  $n$  and  $k$ , and the (equal) spacing of incidence angles  $\Delta\theta$  (in degrees); the program then returns perpendicular polarized, parallel polarized, and unpolarized reflectivities, as well as unpolarized emissivities.

### **totem.f90, totem.cpp, totem.m**

Program `totem` is a routine to evaluate the total, directional or hemispherical emittance or absorptance of an opaque material, based on an array of spectral data, by 10-point Gaussian quadrature.

**Input** (by changing data in the heading of function `emlcl(y)`):

`N` = number of data points for spectral emittance,

`nrefr` = refractive index of adjoining material (`nrefr=1` for vacuum and gases),

`T` = temperature of material (for total emittance), or of gray irradiating source (for total absorptance), in K,

$\lambda(N)$  =  $N$  distinct wavelengths in ascending order, for which the spectral emittance is given, in  $\mu\text{m}$ ,  
 $\text{eps}(N)$  =  $N$  corresponding spectral emittances.

**Output** (printed to screen):

`emitt` = total directional or hemispherical emittance or absorptance.

**Case 1:** Total, directional emittance (`eps` contains spectral, directional values at temperature  $T$ ):

From equation (3.8)

$$\begin{aligned}\epsilon'(T, \hat{\mathbf{s}}) &= \frac{1}{n^2 \sigma T^4} \int_0^\infty \epsilon'_\lambda(\lambda, T, \hat{\mathbf{s}}) E_{b\lambda}(T) d\lambda \\ &= \int_0^1 \epsilon'_\lambda(\lambda(f), T, \hat{\mathbf{s}}) df,\end{aligned}\quad (\text{CC-3-1})$$

where, from equation (1.23)

$$f(n\lambda T) = \int_0^\lambda \frac{E_{b\lambda} d\lambda}{n^2 \sigma T^4}.\quad (\text{CC-3-2})$$

In order to write equation (CC-3-1) in terms of blackbody fraction  $f$ , wavelength must be known as a function of  $f$  (for given  $n$  and  $T$ ), i.e., equation (CC-3-2) must be inverted. The 10 values of  $(n\lambda T)$ , corresponding to the 10 Gaussian quadrature points  $f_i(n\lambda T)$  have been precalculated (using function `bbfn`) and are stored in array `y(i)`. The total emittance is then calculated by expressing equation (CC-3-1) in quadrature form, or

$$\epsilon'(T, \hat{\mathbf{s}}) \simeq \sum_{i=1}^{10} \epsilon'_\lambda(\lambda_i, T, \hat{\mathbf{s}}) w_i,\quad (\text{CC-3-3})$$

where

$$\lambda_i = y_i/nT,\quad (\text{CC-3-4})$$

and the  $w_i$  are Gaussian quadrature weights. This necessitates that  $\epsilon'_\lambda$  must be known at very specific wavelengths, that are not ordinarily part of the given array. The “correct” value for  $\epsilon'_\lambda$  is evaluated by linear interpolation between array values, assuming  $\epsilon'_\lambda = \text{const} = \text{eps}(1)$  for  $\lambda_i < \lambda(1)$ , and  $\epsilon'_\lambda = \text{const} = \text{eps}(N)$  for  $\lambda_i > \lambda(N)$ .

**Case 2:** Total, hemispherical emittance (`eps` contains spectral, hemispherical values at temperature  $T$ ):

From equation (3.10)

$$\begin{aligned}\epsilon(T) &= \frac{1}{n^2 \sigma T^4} \int_0^\infty \epsilon_\lambda(\lambda, T) E_{b\lambda} d\lambda = \int_0^1 \epsilon_\lambda(\lambda(f), T) df \\ &\simeq \sum_{i=1}^{10} \epsilon_\lambda(\lambda_i, T) w_i.\end{aligned}\quad (\text{CC-3-5})$$

Thus, the calculation is identical to Case 1.

**Case 3:** Total, directional absorptance (`eps` contains spectral, directional values at the surface temperature  $T_s$ , irradiation is assumed to come from a gray source at temperature  $T$ ).

From equations (3.23) and (3.31)

$$\begin{aligned}\alpha'(T_s, T, \hat{\mathbf{s}}) &= \frac{1}{n^2 \sigma T^4} \int_0^\infty \epsilon'_\lambda(\lambda, T, \hat{\mathbf{s}}) E_{b\lambda}(T) d\lambda \\ &= \int_0^1 \epsilon'_\lambda(\lambda(f), T_s) df \simeq \sum_{i=1}^{10} \epsilon'_\lambda(\lambda_i, T_s) w_i,\end{aligned}\quad (\text{CC-3-6})$$

and the calculation is again identical.

**Case 4:** Total, hemispherical absorptance (`eps` contains spectral, hemispherical values at surface temperature  $T_s$ ; irradiation is assumed to be gray and diffuse with source temperature  $T$ ).

Then, from equations (3.27) and (3.31)

$$\begin{aligned}\alpha(T_s, T) &= \frac{1}{n^2 \sigma T^4} \int_0^\infty \epsilon_\lambda(\lambda, T_s) E_{b\lambda}(T) d\lambda \\ &= \int_0^1 \epsilon_\lambda(\lambda(f), T_s) df \approx \sum_{i=1}^{10} \epsilon_\lambda(\lambda_i, T_s) w_i.\end{aligned}\tag{CC-3-7}$$

### Examples

Two examples have been programmed into `totem` (or, rather, function `emlcl`):

1.: The material of Problem 3.1, with a step function in spectral emittance of

$$\epsilon_\lambda = \begin{cases} 0.5, & \lambda < 5\mu\text{m}, \\ 0.3, & \lambda > 5\mu\text{m}, \end{cases}$$

and a temperature of  $T = 500$  K. For part *a*) `nrefr=1.0`, and for *b*) `nrefr=2.0` (implemented here) This results in `emitt=0.3435` for *a*) and `emitt=0.4296` for *b*).

2.: Aluminum oxide, as given in Fig. 1-14, discretized into eight equally-spaced values (commented out as given here). For temperature of  $T = 500$  K and `nrefr=1.0` this results in `emitt=0.7494`.